

BU

Solution to Assignment 1.

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1 (a). The minus sign means: If $\frac{\partial \phi}{\partial x} > 0$ for $a \leq x \leq b$, ϕ is an increasing function of x . The heat energy in this interval must decrease ($\frac{\partial Q}{\partial t} < 0$).

(2) (b) The minus sign means: If the temperature increases as x increases ($\frac{\partial \phi}{\partial x} > 0$), then the heat energy flows to the left.

2. (a) The total energy between x and $x+\Delta x$ is

$$\int_x^{x+\Delta x} e \cdot A dx$$

(4) Since $Q=0$, we have

$$\frac{d}{dt} \int_x^{x+\Delta x} e \cdot A dx = \phi(x,t)A - \phi(x+\Delta x,t)A$$

when Δx is small, $\int_x^{x+\Delta x} e \cdot A dx \approx e \cdot A \cdot \Delta x$

So we have

$$\frac{\partial}{\partial t} (e(x,t) \cdot A \cdot \Delta x) = \phi(x,t)A - \phi(x+\Delta x,t)A$$

let $\Delta x \rightarrow 0$:

$$\frac{\partial e}{\partial t} = \lim_{\Delta x \rightarrow 0} \frac{\phi(x,t) - \phi(x+\Delta x,t)}{\Delta x} = -\frac{\partial \phi}{\partial x}$$

(b) $[a, b]$ The total energy is

$$\int_a^b e \cdot A dx$$

The conservation law gives

$$\frac{d}{dt} \int_a^b e \cdot A dx = A \phi(a,t) - \phi(b,t) \cdot A = A \int_a^b -\frac{\partial \phi}{\partial x} dx$$

$$\int_a^b \left(\frac{\partial e}{\partial t} + \frac{\partial \phi}{\partial x} \right) dx = 0$$

3. The latter condition at $x=x_0$ is.

$$k_0(x_{0+}) \cdot \frac{\partial u}{\partial x}(x_{0+}, t) = k_0(x_{0-}) \frac{\partial u}{\partial x}(x_{0-}, t).$$

(4) In the case $k_0(x_{0+}) = k_0(x_{0-})$, we have

$$\frac{\partial u}{\partial x}(x_{0+}, t) = \frac{\partial u}{\partial x}(x_{0-}, t).$$

So in this case, $\frac{\partial u}{\partial x}$ is continuous at $x=x_0$.

4. (a) $\rho c \frac{\partial u}{\partial t} = k_0 \frac{\partial^2 u}{\partial x^2} + Q, \quad Q=0$

$$\left\{ \begin{array}{l} u(0) = T, \quad u(L) = 0 \end{array} \right.$$

(5) equilibrium: $\left\{ \begin{array}{l} \frac{d^2 u}{dx^2} \equiv 0 \\ (\frac{\partial u}{\partial t} \equiv 0) \\ u(0) = T, \quad u(L) = 0 \end{array} \right.$

$$u = c_1 x + c_2; \quad \begin{array}{l} u(0) = T \Leftrightarrow c_2 = T \\ u(L) = 0 \Leftrightarrow c_1 = -\frac{T}{L} \end{array}$$

$$u(x) = -\frac{T}{L} x + T.$$

(b) $\frac{\partial u}{\partial t} = 0, \quad \frac{Q}{k_0} = 1.$

equilibrium $\left\{ \begin{array}{l} \frac{d^2 u}{dx^2} + 1 = 0 \\ u(0) = T_1, \quad u(L) = T_2. \end{array} \right.$

(6) $u = -\frac{x^2}{2} + c_1 x + c_2$

$$u(0) = T_1 \Rightarrow c_2 = T_1$$

$$u(L) = T_2 \Rightarrow -\frac{L^2}{2} + c_1 L + T_1 = T_2 \Rightarrow c_1 = \frac{1}{L} \left(T_2 - T_1 + \frac{L^2}{2} \right)$$

$$\therefore u(x) = -\frac{x^2}{2} + \frac{1}{L} \left(T_2 - T_1 + \frac{L^2}{2} \right) x + T_1$$

(c) $u_{xx} = 0 \Rightarrow u = c_1 x + c_2, \quad u(0) = T \Rightarrow c_2 = T.$

(7) $u_x(L) + u(L) = 0 \Rightarrow c_1 + c_1 L + c_2 = 0 \Rightarrow c_1 = -\frac{T}{c_1 + L}.$

$$u = -\frac{T}{c_1 + L} x + T$$

5 Equation for $0 < x < 1$

$$cp \frac{\partial u}{\partial t} = k_0 \frac{\partial^2 u}{\partial x^2} + Q$$

$$\begin{cases} u(0) = 0 \\ cp = 1, Q = 1 \\ k_0 = 1 \end{cases}$$

Equation for $1 < x < 2$

$$cp \frac{\partial u}{\partial t} = k_0 \frac{\partial^2 u}{\partial x^2} + Q$$

$$\begin{cases} u(2) = 0 \\ Q = 0, cp = 1, k_0 = 2 \end{cases}$$

at $x = 1$ we have

$$\begin{cases} u(1^-) = u(1^+) \\ k_0(1^-) \frac{\partial u(1^-)}{\partial x} = k_0(1^+) \frac{\partial u(1^+)}{\partial x} \end{cases}$$

Equilibrium

when $0 < x < 1$:

$$\frac{d^2 u}{dx^2} + 1 = 0$$

$$u = -\frac{x^2}{2} + c_1 x + c_2$$

$$u(0) = 0$$

$$u(0) = 0 \Rightarrow c_2 = 0$$

$$u = -\frac{x^2}{2} + c_1 x, \quad 0 < x < 1$$

Equilibrium when $1 < x < 2$:

$$\frac{d^2 u}{dx^2} = 0$$

$$u = c_3 + c_4 x$$

$$u(2) = 0$$

$$\Rightarrow c_3 = -2c_4$$

$$u = c_4(-2 + x)$$

At $x = 1$:

$$-\frac{x^2}{2} + c_1 x \Big|_{x=1^-} = c_4(-2+x) \Big|_{x=1^+}$$

$$1 \cdot \frac{\partial u}{\partial x} \Big|_{x=1^-} = 2 \cdot \frac{\partial u}{\partial x} \Big|_{x=1^+}$$

i.e.

$$-\frac{1}{2} + c_1 = -c_4$$

$$1 \cdot (-x + c_1) \Big|_{x=1^-} = 2 \cdot c_4$$

$$c_4 = \frac{1}{6}, \quad c_1 = \frac{2}{3}$$

$$u = \frac{1}{6}(-2+x), \quad 1 < x < 2$$

and

$$u = -\frac{x^2}{2} + \frac{2}{3}x, \quad 0 < x < 1$$

6 (a) Equilibrium equation

$$\begin{cases} k_0 u_{xx} + Q_0 = 0 \\ u_x(0) = 0, u_x(L) = 0 \end{cases} \Rightarrow u = \frac{-Q_0}{2k_0} x^2 + c_1 x + c_2$$

$$u_x(0) = 0 \Rightarrow c_1 = 0$$

$$u_x(L) = 0 \Rightarrow -\frac{Q_0}{k_0} L + c_1 = 0$$

So it require $Q_0 L/k_0 = 0$.

this is impossible. No existence of equilibrium temperature distribution.

physically: since two ends of the rod are insulated, and there are a source term Q_0 inside the rod, the temperature of the rod will tend to infinity if $Q_0 > 0$.

$$(b) \begin{cases} \rho c \frac{\partial u}{\partial t} = k_0 u_{xx} + Q_0 \\ u_x(0) = u_x(L) = 0 \end{cases}$$

Total energy $E = A \int_0^L \rho c \cdot u(x,t) dx$.

$$\frac{dE}{dt} = A \int_0^L \rho c \cdot u_t(x,t) dx$$

$$= A \int_0^L \rho c \cdot (k_0 u_{xx} + Q_0) dx \quad (\text{using the equation})$$

$$= 0 + A \int_0^L \rho c Q_0 dx \quad (\text{using the bc})$$

$$\therefore \frac{dE}{dt} = A \int_0^L \rho c Q_0 dx = \begin{cases} > 0, & \text{if } Q_0 > 0 \\ < 0, & \text{if } Q_0 < 0 \end{cases}$$

$$E(t) = E(0) + A \cdot t \int_0^L \rho c Q_0 dx$$

E is an increasing function if $Q_0 > 0$ and an decreasing function if $Q_0 < 0$.

$$7. (a) \begin{cases} u_{xy} = 0 & \Rightarrow u = c_1 x + c_2 \\ u_x(0) = 1 & u_x(L) = \beta \end{cases}$$

$$\textcircled{4} \quad \begin{aligned} u_x(0) = 1 & \Rightarrow c_1 = 1 \\ u_x(L) = \beta & \Rightarrow c_1 = \beta \end{aligned}$$

When $\beta = 1$, there are solutions.

$$(b) \begin{cases} u'' + x - \beta = 0 & u = -\frac{1}{6}(x-\beta)^3 + c_1 x + c_2 \\ u_x(0) = 0 & u'(L) = 0 \end{cases}$$

$$\textcircled{4} \quad \begin{aligned} u_x(0) = 0 & \Rightarrow -\frac{1}{2}(x-\beta)^2 \Big|_{x=0} + c_1 = 0 \Rightarrow -\frac{1}{2}\beta^2 + c_1 = 0 \\ u'(L) = 0 & \Rightarrow -\frac{1}{2}(L-\beta)^2 + c_1 = 0 \Rightarrow -\frac{1}{2}(L-\beta)^2 + c_1 = 0 \end{aligned}$$

$$\therefore \beta^2 = (L-\beta)^2 \Rightarrow \beta = \frac{L}{2}$$

When $\beta = \frac{L}{2}$, there are solutions.

$$8 \quad \frac{\partial u}{\partial x} = \frac{\partial u}{\partial r} \frac{\partial r}{\partial x} + \frac{\partial u}{\partial \theta} \frac{\partial \theta}{\partial x} = u_r \cdot (\sqrt{x^2+y^2})_x + u_\theta \cdot (\arctan \frac{y}{x})_x$$

$$= u_r \cdot \frac{x}{(x^2+y^2)^{\frac{1}{2}}} + u_\theta \cdot \frac{-y}{x^2+y^2}$$

$$\textcircled{4} \quad u_y = u_r \cdot \frac{y}{(x^2+y^2)^{\frac{1}{2}}} + u_\theta \cdot \frac{x}{x^2+y^2}$$

Continue the process to have $\frac{\partial^2 u}{\partial x^2} + \frac{\partial^2 u}{\partial y^2} = \frac{\partial^2 u}{\partial r^2} + \frac{1}{r} \frac{\partial u}{\partial r} + \frac{1}{r^2} \frac{\partial^2 u}{\partial \theta^2}$

Method 2: $\frac{\partial u}{\partial r} = u_x \cos \theta + u_y \sin \theta$

$$\frac{\partial u}{\partial \theta} = u_x (-r \sin \theta) + u_y (r \cos \theta)$$

Finally $\frac{1}{r} \frac{\partial}{\partial r} (r \frac{\partial u}{\partial r}) + \frac{1}{r^2} \frac{\partial^2 u}{\partial \theta^2} = (\cos^2 \theta + \sin^2 \theta) u_{xx} + (\cos^2 \theta + \sin^2 \theta) u_{yy}$
 $= u_{xx} + u_{yy}$