

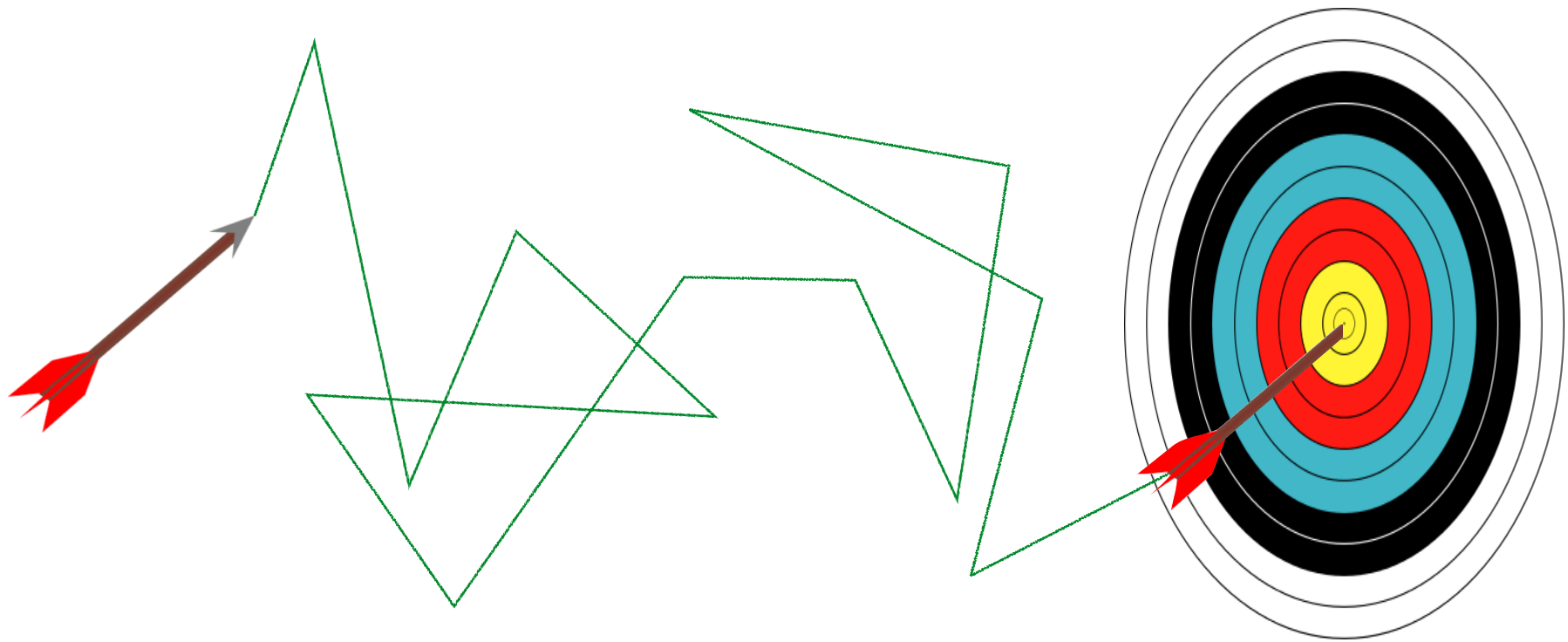
# Problems in diffusion and absorption: How fast can you hit a target with a random walk?

Andrew J. Bernoff  
Harvey Mudd College

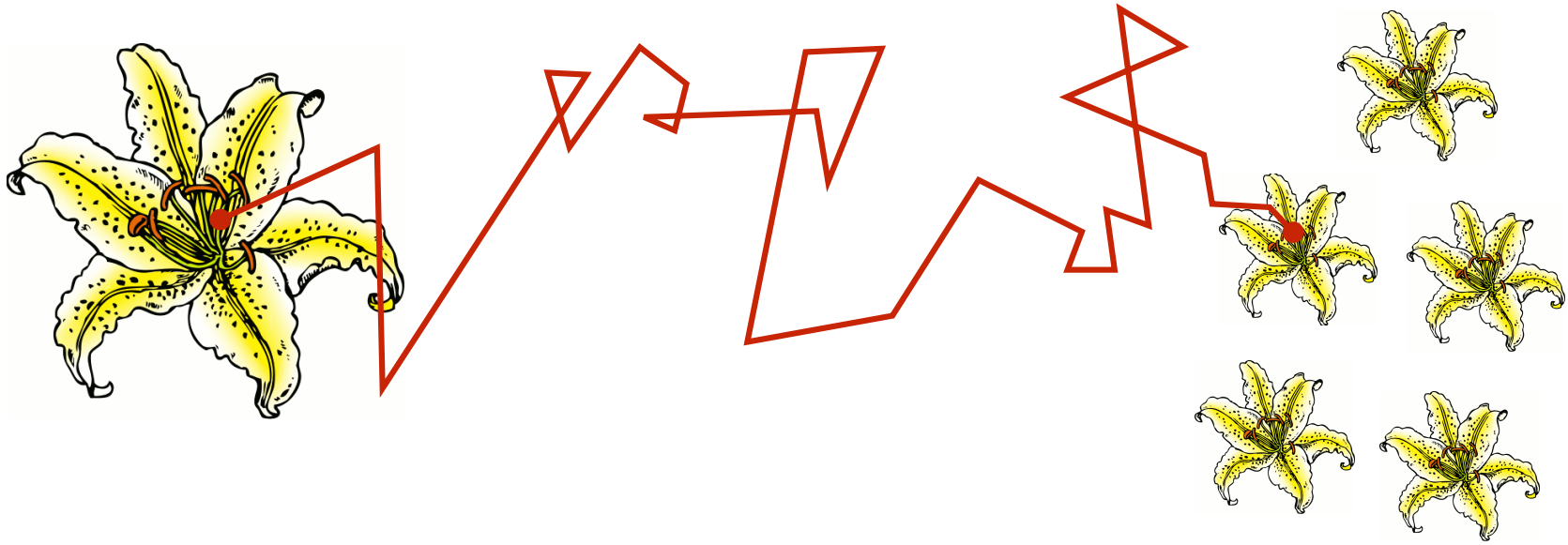
In collaboration with **Alan Lindsay** (Notre Dame)

Thanks to Alan Lindsay, Michael Ward (UBC), Justin Tzou (UBC) and Theodore Kolokolnikov for introducing me to these problems

# Random Signaling Problems

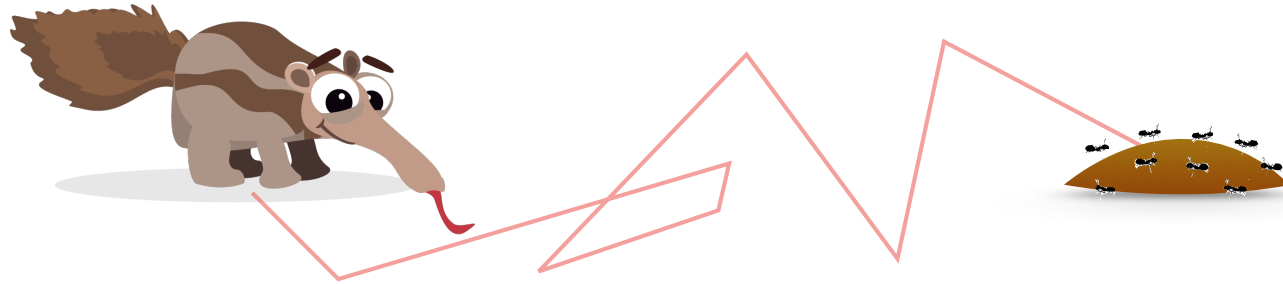


# Pollination



- *What what probability will a pollen grain released from the stamen of one flower find the pistil of another flower?*
- *On average, how long does this take?*
- *How does the shape of the stamen affect this?*

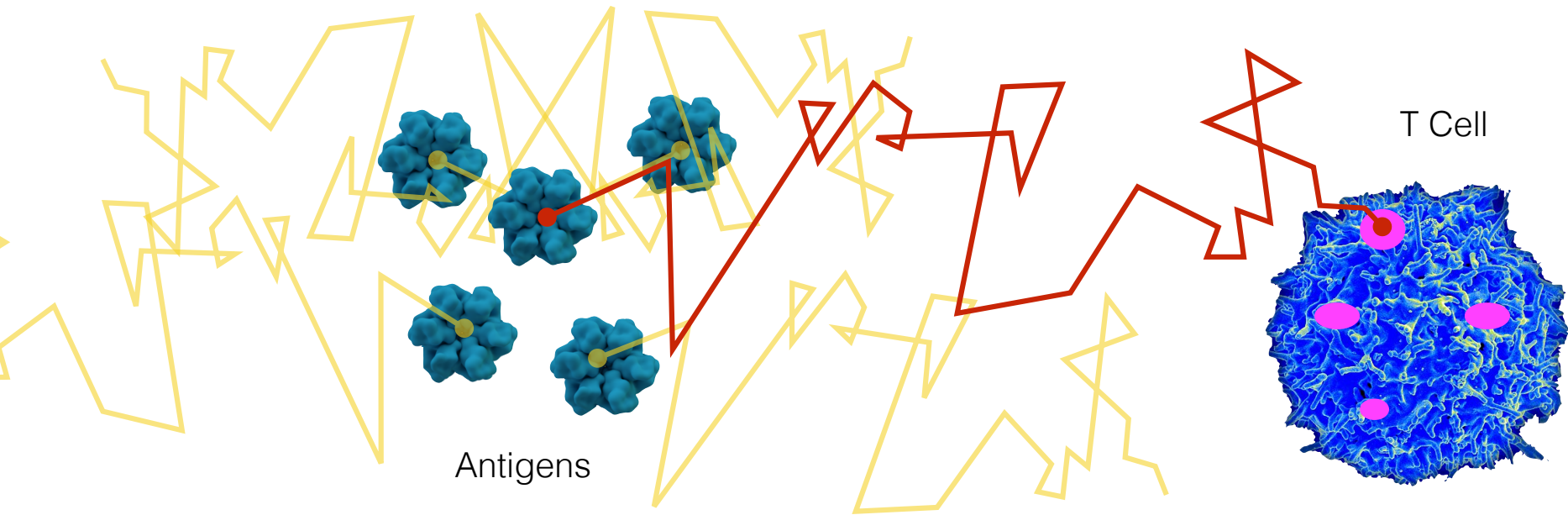
# Foraging



Suppose an animal is randomly foraging for food.

- *Will the anteater find the ants?*
- *On average, how long does this take?*

# Molecular Signaling



When an antigen (a toxin or a protein that promotes an immune response) binds to a receptor on a T-cell it can trigger the creation of antibodies.

- *What is the probability of this binding occurring?*
- *On average, how long does this take?*
- *How does the distribution of the receptors affect this?*

# Modeling Diffusion & Capture



pollen on a stamen

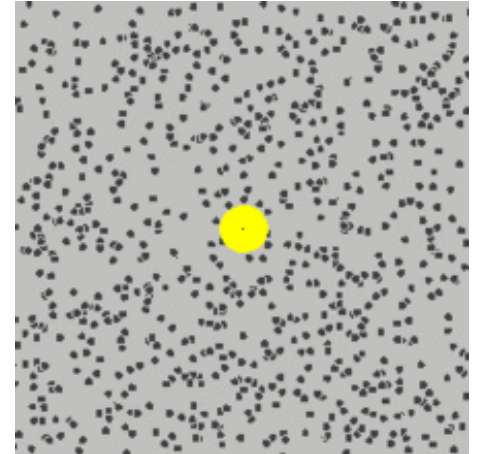
# Brownian Motion

Brownian motion is the random walk a molecule takes due to collisions with other molecules.

Typically it can be describe by a distribution of velocities.

Let

- $V_0$  be the average speed
- $L_0$  be the mean free path



[https://en.wikipedia.org/wiki/Brownian\\_motion](https://en.wikipedia.org/wiki/Brownian_motion)

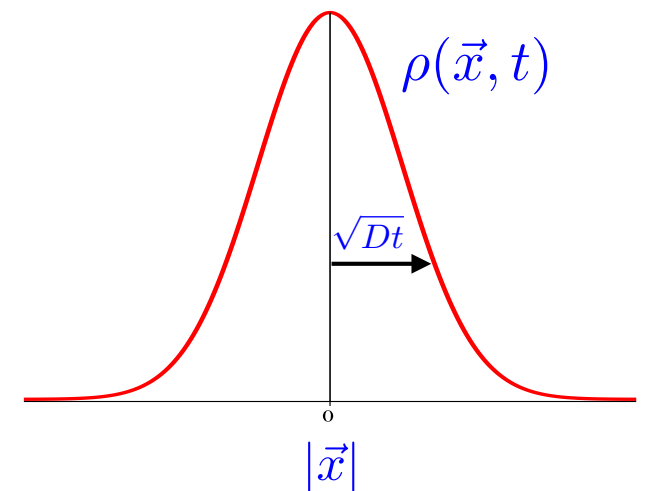
# Random Walks

Due to the **Central Limit Theorem** after many steps the density,  $\rho(\vec{x}, t)$ , of the random walk converges to a Gaussian distribution.

In N-dimensions, the distribution approaches

$$\rho(\vec{x}, t) = \left( \frac{1}{\sqrt{4\pi Dt}} \right)^N e^{-|\vec{x}|^2 / (4Dt)}$$

$$\vec{x} \in \mathbb{R}^N$$



The diffusion constant,  $D$ , scales as

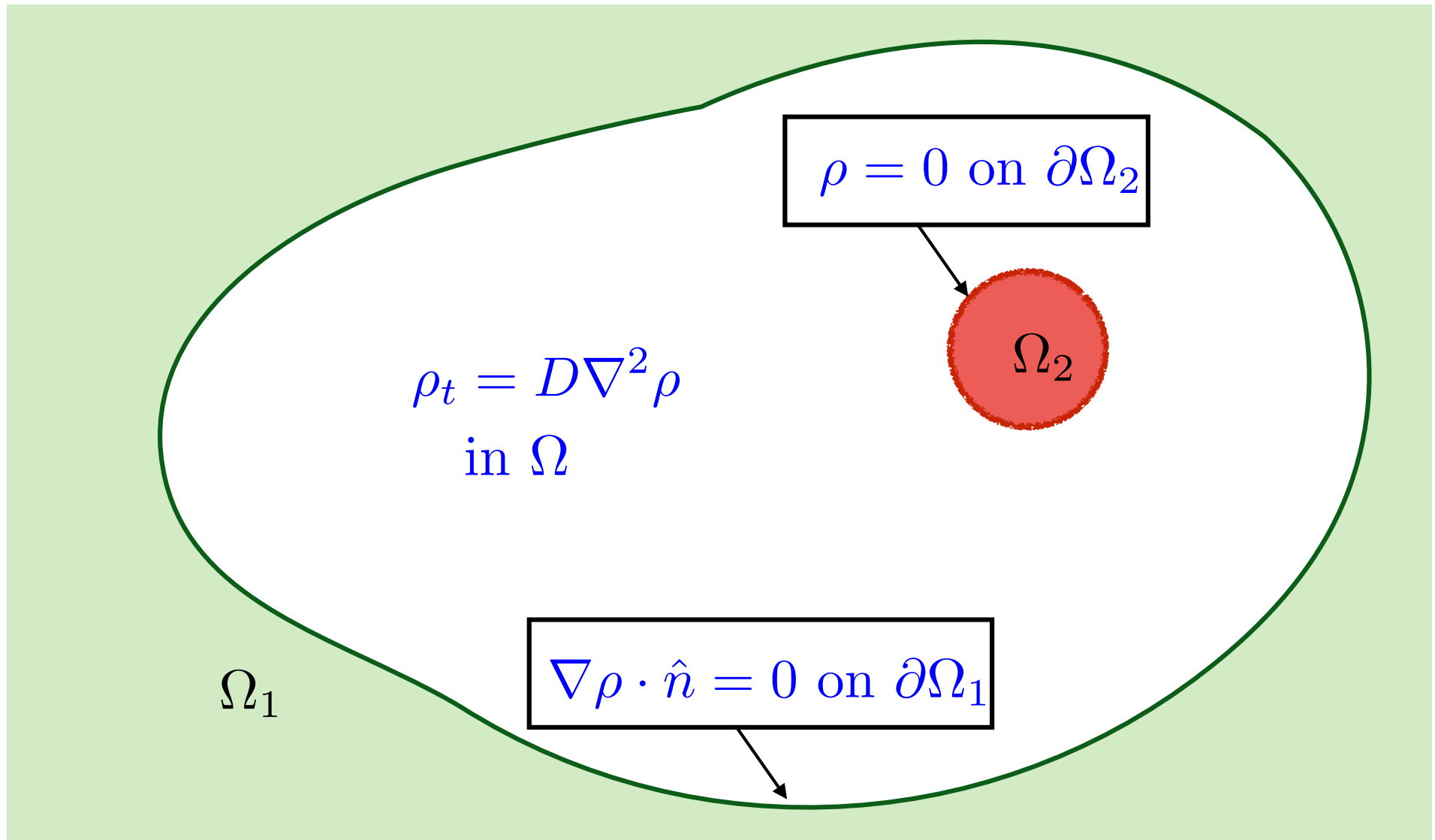
$$D \propto (V_0)^2 / L_0 \quad [\text{Length}]^2 \times [\text{Time}]$$

We call  $\rho(\vec{x}, t)$  the *diffusion kernel*.



# The Diffusion Equation

The density of molecules,  $\rho(\vec{x}, t)$ , satisfies the diffusion equation in some domain  $\Omega$ .



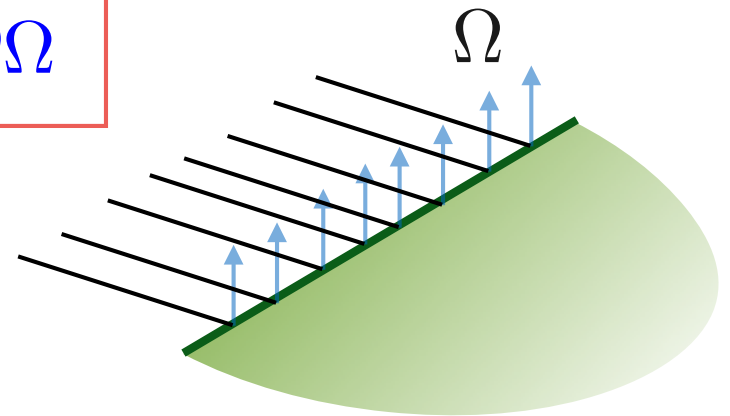
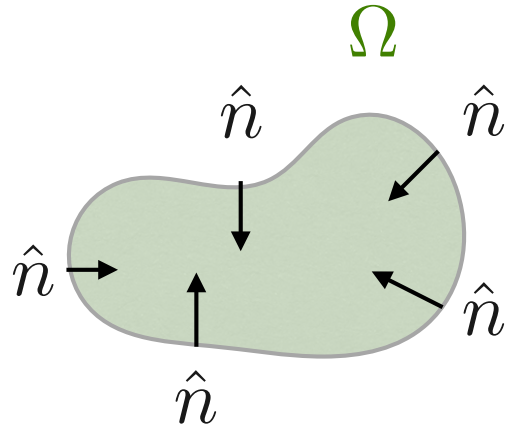
# Flux & Neumann BC's

The flux of molecules is given by the Fourier Law

$$\vec{F} = D\nabla u \quad [\text{velocity}] \times [\text{density}]$$

At an impermeable boundary, we have a no-flux BC:

$$\vec{F} \cdot \hat{n} = 0 \quad \Rightarrow \quad \nabla u \cdot \hat{n} = 0 \text{ on } \partial\Omega$$

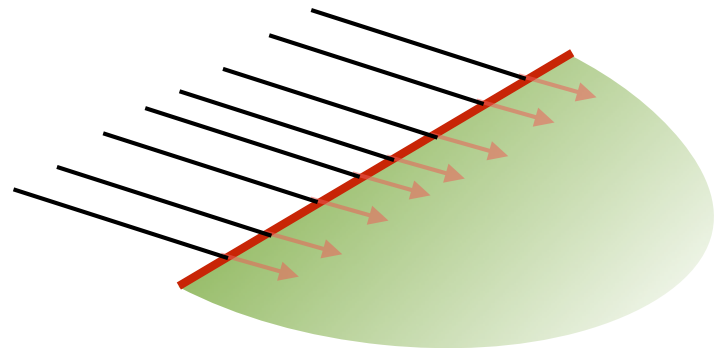
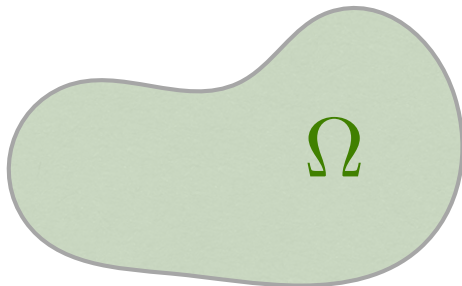


At the particle level, particles are reflected.  
Angle of incidence = Angle of reflection

# Capture & Dirichlet BC's

At a boundary that captures every molecule that impacts it the density will drop to zero

$$u = 0 \text{ on } \partial\Omega$$



At the particle level, particles are captured.

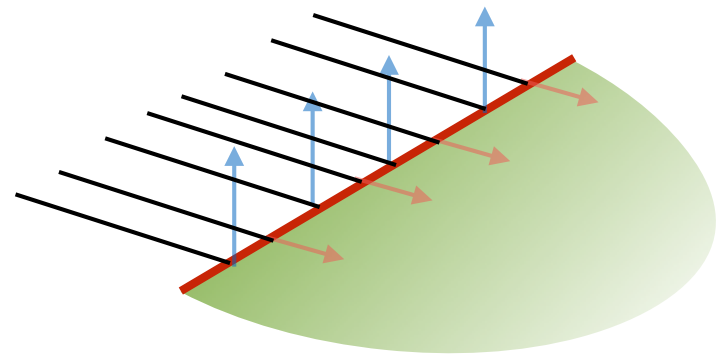
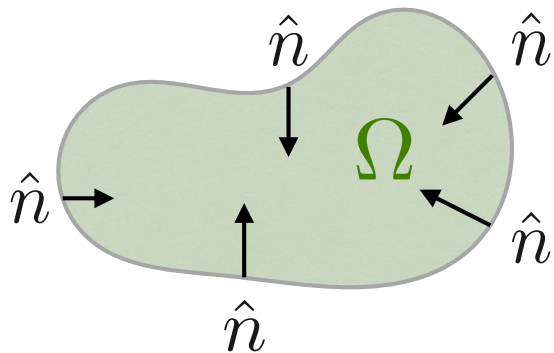
# Partial Capture & Robin BC's

Consider a boundary that reflects some, but not all, of the particles that strike it.

$$\nabla u \cdot \hat{n} + \kappa u = 0 \text{ on } \partial\Omega$$

$\kappa \ll 1$  Mostly Reflecting

$\kappa \gg 1$  Mostly Absorbing



At the particle level, some particles are captured while some are reflected.

# Metrics for Capture



# Probability of Capture

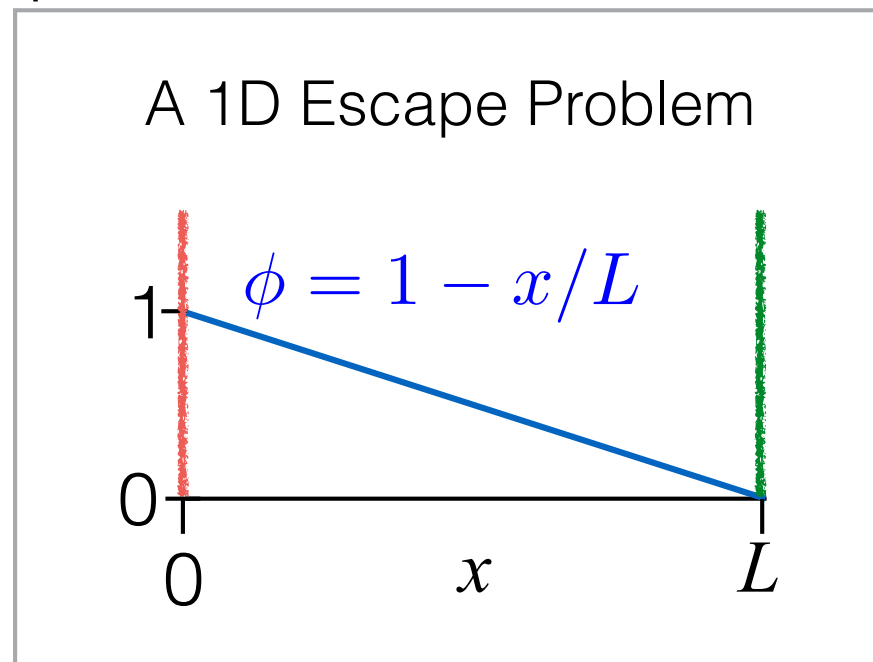
Consider a domain with a target boundary and an escape boundary.

The probability,  $\phi(x)$ , that a molecule is captured satisfies Laplace's equation:

$$\nabla^2 \phi = 0 \text{ in } \Omega$$

$$\phi = 1 \text{ on } \partial\Omega \text{ for capture}$$

$$\phi = 0 \text{ on } \partial\Omega \text{ for escape}$$



This is sometimes called *harmonic measure*.

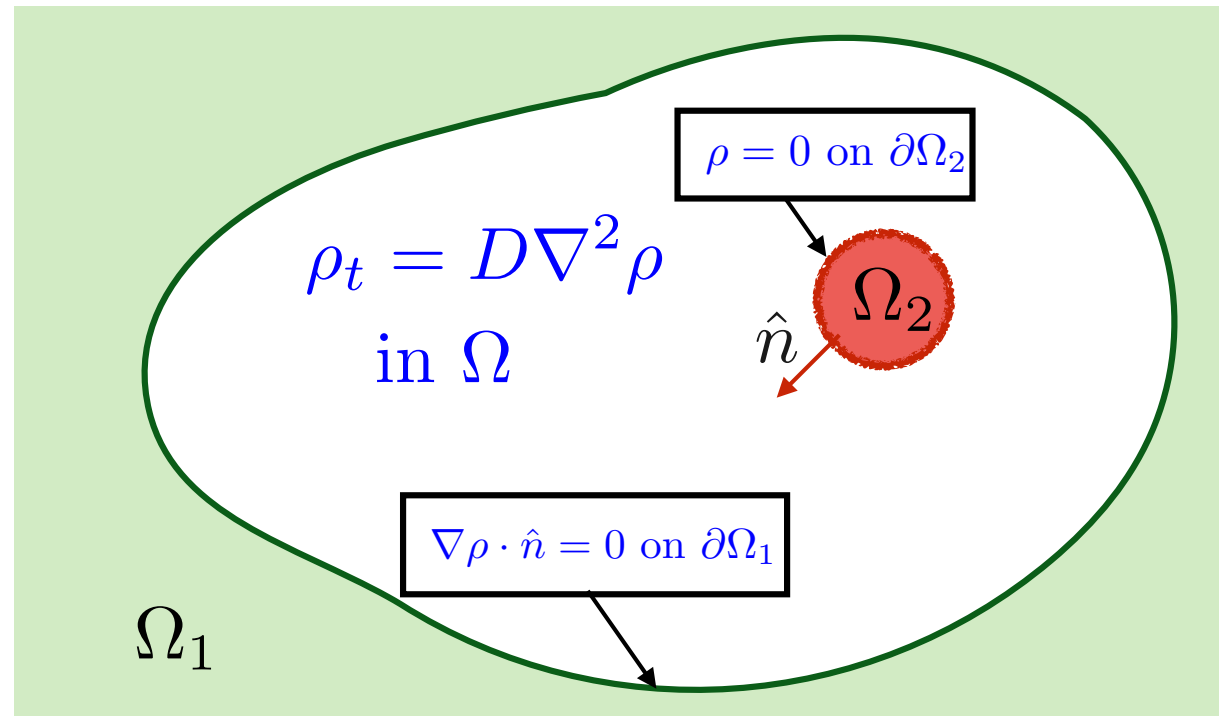
# Distribution of First Passage Time (FPT)

Consider a domain with a target boundary.

The density of mass capture per unit time,  $p(t)$ , can be computed by integrating the flux over the target boundary.

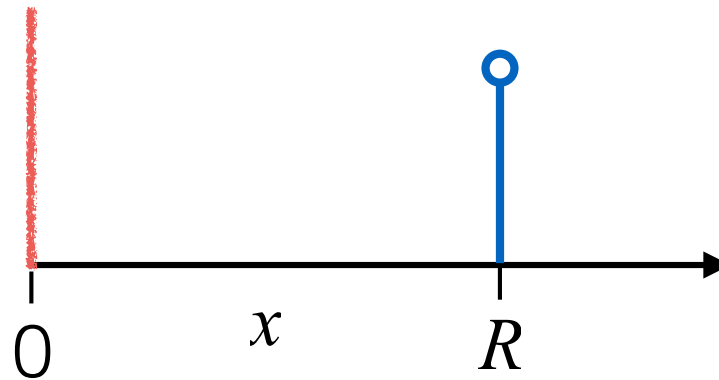
$$p(t) = - \oint_{\Omega_2} D \nabla \rho \cdot \hat{n} dS$$

$$\frac{[\text{Mass}]}{[\text{Time}]}$$

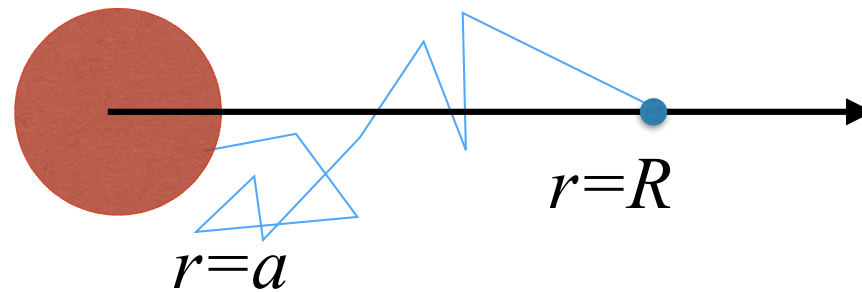


# The Importance of Dimension

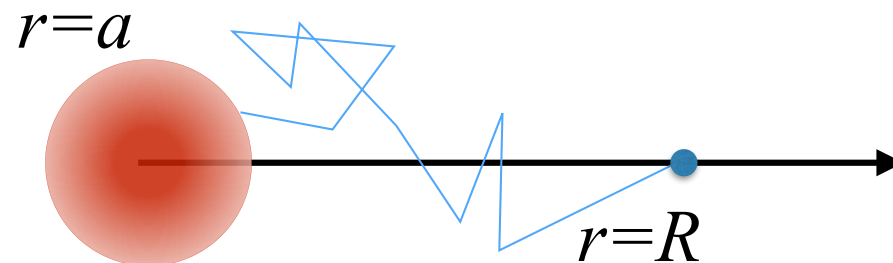
1D



2D



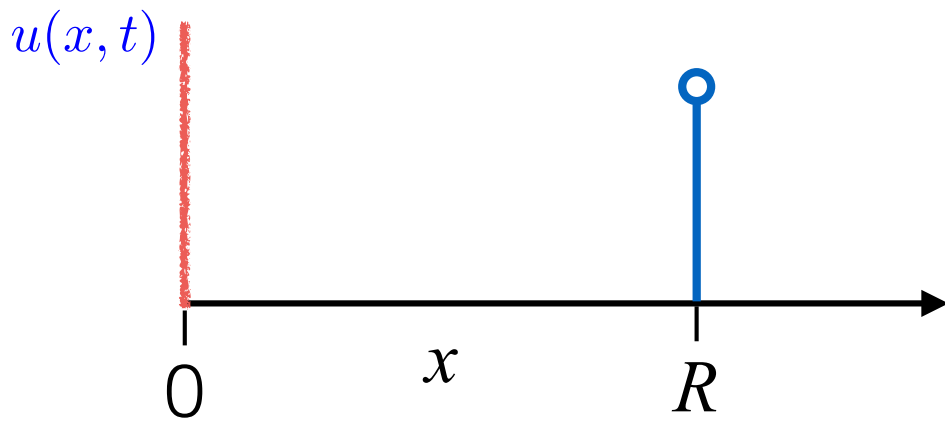
3D





# Capture in 1D

Consider a  $\delta$ -function release at  $x=R$ .



$$u_t = Du_{xx} \quad x > 0, t > 0$$

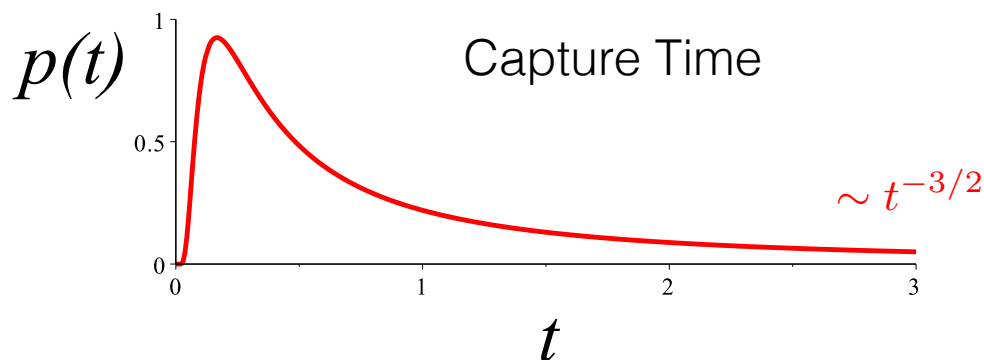
$$u(x, 0) = \delta(x - R) \quad x > 0$$

$$u(0, t) = 0 \quad t > 0$$

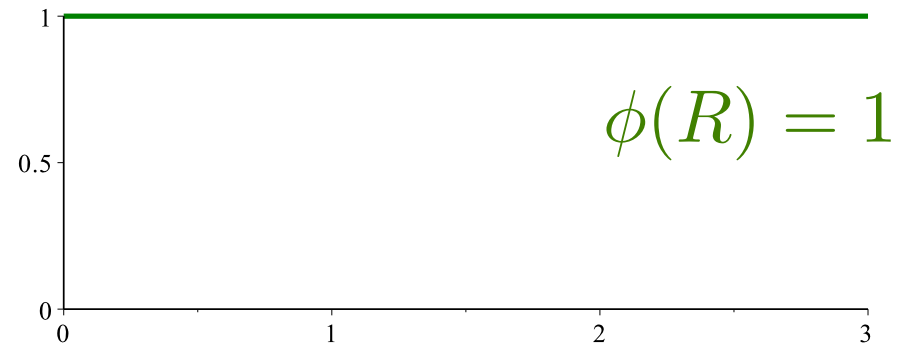
By method of images:

$$u(x, t) = \frac{1}{2\sqrt{\pi Dt}} [e^{-(x-R)^2/4Dt} - e^{-(x+R)^2/4Dt}]$$

FPT:  $p(t) = -Du_x(0, t) = \frac{Re^{-R^2/4Dt}}{2\sqrt{\pi Dt^{3/2}}}$

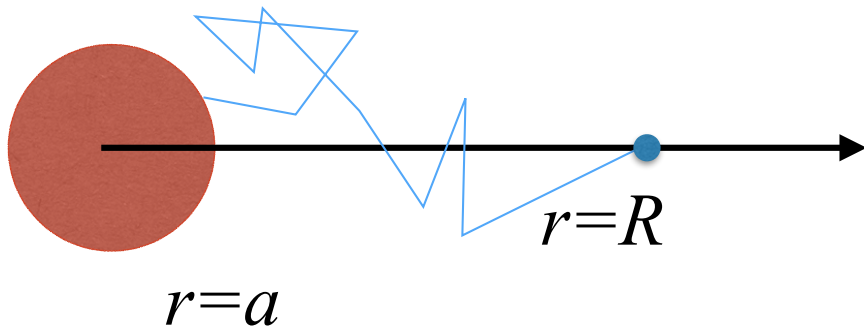


Probability of capture:



# Capture in 2D

Consider a  $\delta$ -function release at  $r=R$  and a circular trap of radius  $a$ .



$$u_t = D \left[ u_{rr} + \frac{1}{r} u_r \right] \quad r > 0, t > 0$$

$$u(r, 0) = \frac{1}{2\pi R} \delta(r - R) \quad r > a$$

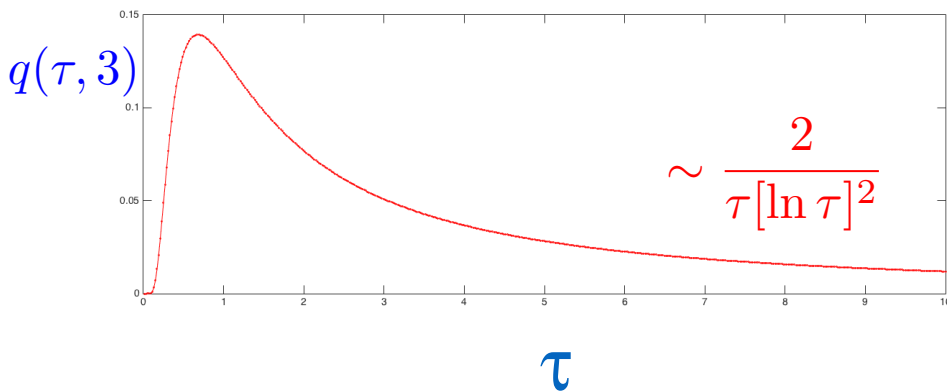
$$u(a, t) = 0 \quad t > 0$$

FPT:  $p(t) = \frac{a^2}{d} q(\tau, \delta) \quad \delta = \frac{R}{a} \quad \tau = \frac{a^2 t}{d}$

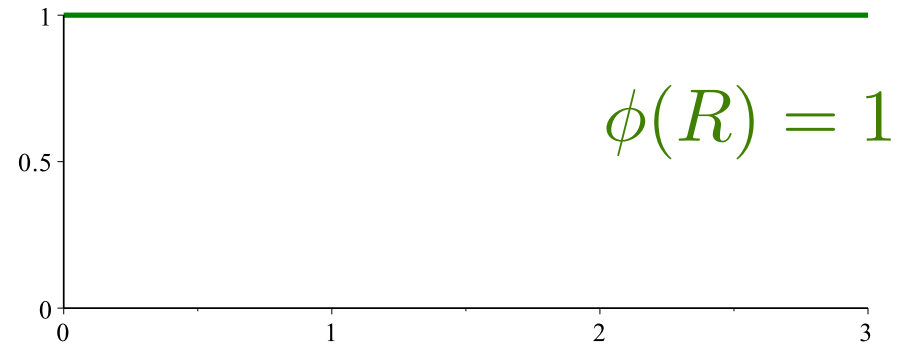
Obtained with a  
Laplace Transform

$$q(\tau, \delta) = \frac{2}{\pi} \int_{s=0}^{\infty} s e^{-s^2 \tau} \frac{J_0(s) Y_0(\delta s) - J_0(\delta s) Y_0(s)}{[J_0(s)]^2 + [Y_0(s)]^2} ds$$

Capture Time

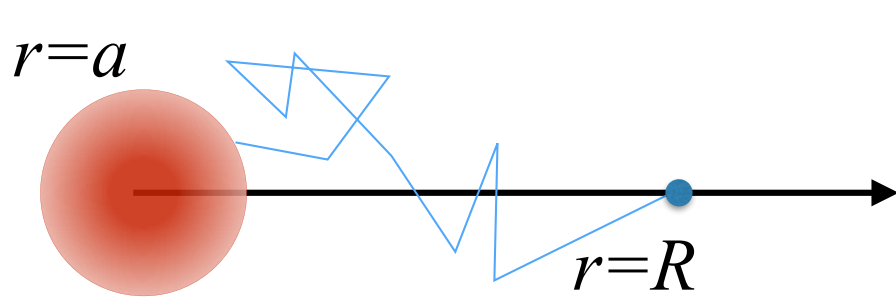


Probability of capture:



# Capture in 3D

Consider a  $\delta$ -function release at  $r=R$  and a spherical trap of radius  $a$ .



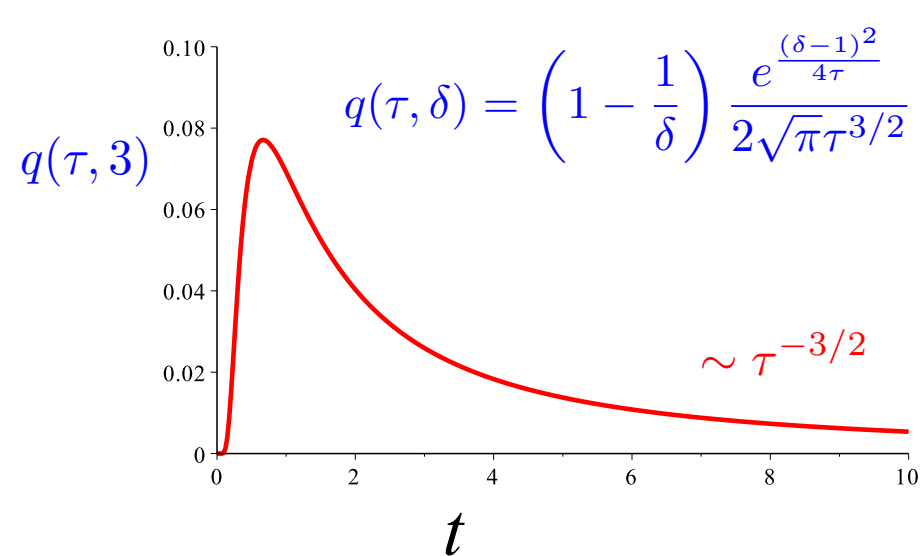
$$u_t = D \left[ u_{rr} + \frac{2}{r} u_r \right] \quad r > 0, t > 0$$

$$u(r, 0) = \frac{1}{4\pi R^2} \delta(r - R) \quad r > a$$

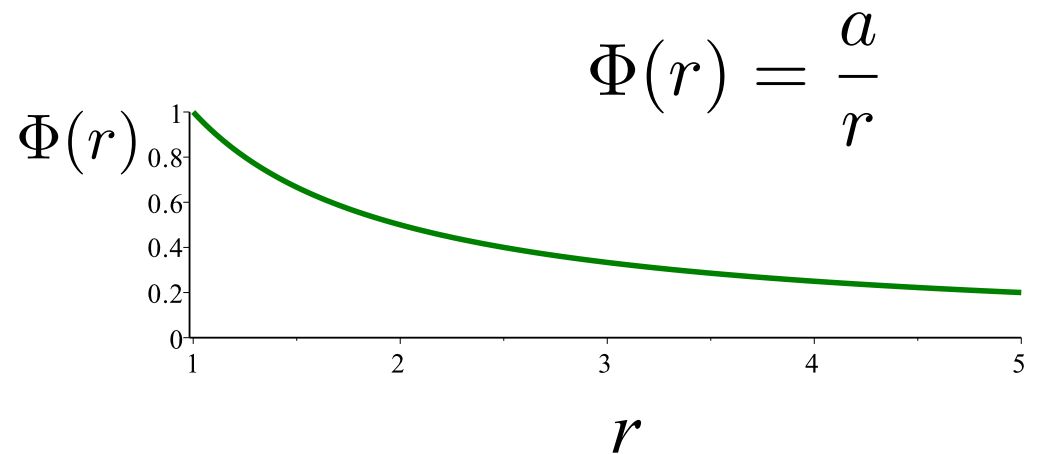
$$u(a, t) = 0 \quad t > 0$$

FPT:  $p(t) = \frac{a^2}{d} q(\tau, \delta) \quad \delta = \frac{R}{a} \quad \tau = \frac{a^2 t}{d}$

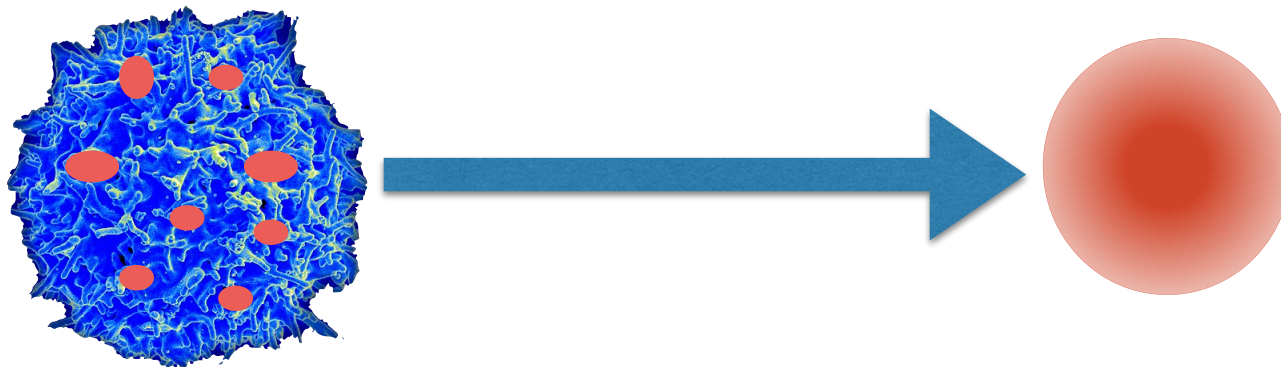
Capture Time



Probability of capture:

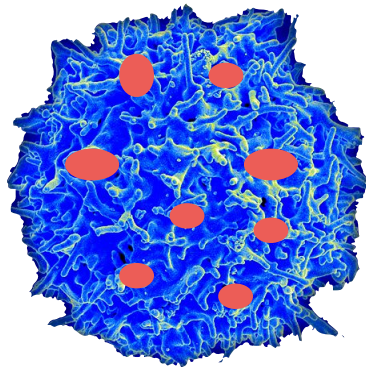


# Homogenization & Effective Boundary Conditions



# The Basic Question

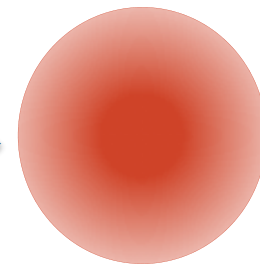
Can you replace a complicated boundary condition or geometry with a much simpler one?



Irregular shape cell  
Many pores

Asymptotics

Numerics

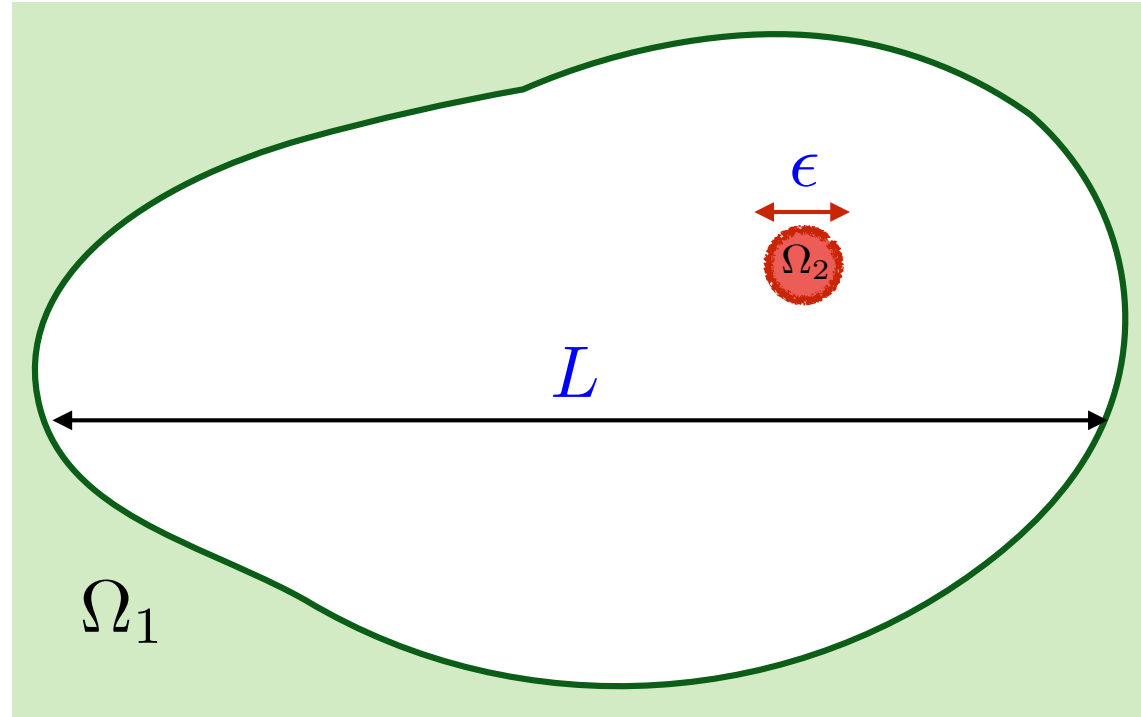


Spherical Trap

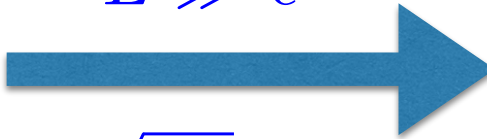
# The steady-state ansatz

Suppose we are looking at a **small** target in a **large** domain.

When can we assume that the density is steady near the target?



$$\begin{aligned}\rho_t &= D\nabla^2\rho \text{ in } \Omega \\ \rho &= 0 \text{ on } \partial\Omega_2 \\ \nabla\rho \cdot \hat{n} &= 0 \text{ on } \partial\Omega_1\end{aligned}$$

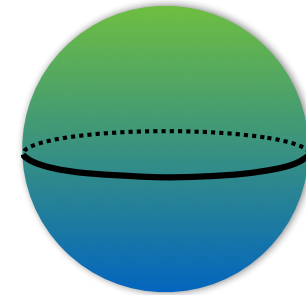
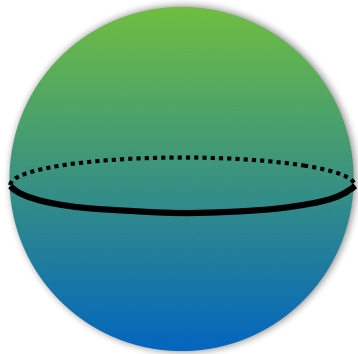
$$\begin{aligned}L &\gg \epsilon \\ \sqrt{Dt} &\gg \epsilon\end{aligned}$$


$$\begin{aligned}\nabla^2\rho &= 0 \text{ in } \Omega \\ \rho &= 0 \text{ on } \partial\Omega_2 \\ \rho &\rightarrow \rho_0 \text{ for } \epsilon \ll |\vec{x}| \ll L\end{aligned}$$

# A simple example

Sphere with a partial absorbing boundary

Smaller sphere with a totally absorbing boundary



$$\nabla^2 u = 0 \quad r > a$$

$$u_r(a) - \kappa u(a) = 0$$

$$\lim_{r \rightarrow \infty} u(r) = C_\infty$$

$$\nabla^2 u = 0 \quad r > b$$

$$u(b) = 0$$

$$\lim_{r \rightarrow \infty} u(r) = C_\infty$$

$$u(r) = C_\infty \left( 1 - \frac{1}{1 + 1/\kappa a} \frac{a}{r} \right)$$

$$u(r) = C_\infty \left( 1 - \frac{b}{r} \right)$$

Problems are equivalent when 
$$b = \frac{a}{1 + 1/\kappa a}$$

# A Pore on a Plane

Consider a circular pore of radius  $a$  on an infinite plane.

There is an exact solution for  $u(r, z)$ .

$$\nabla^2 u = 0 \quad z > 0$$

$$u_z(r, 0) = 0 \text{ for } r > a$$

$$u(r, 0) = 0 \text{ for } r < a$$

$$\lim_{z \rightarrow \infty} u(r, z) = C_\infty$$

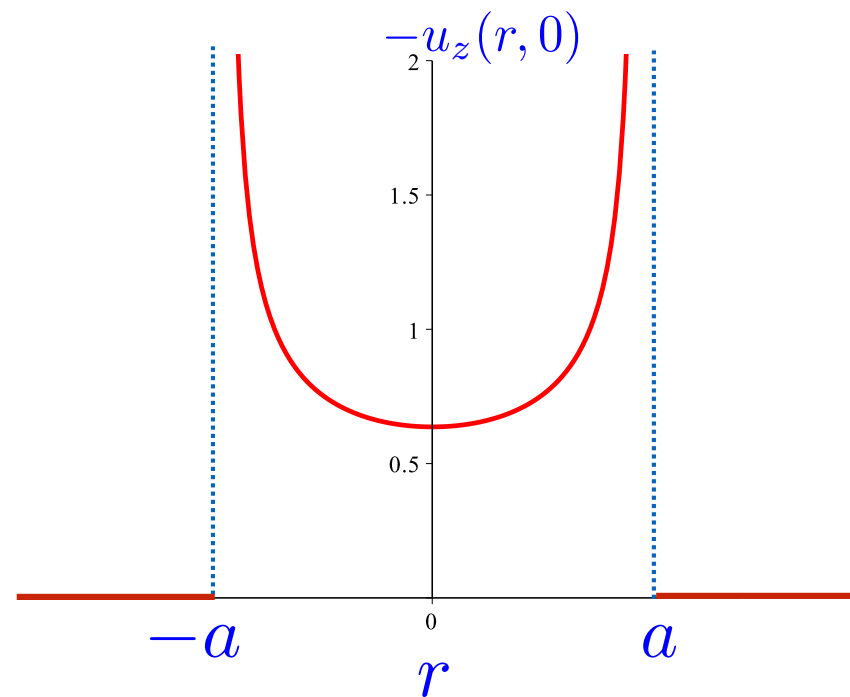
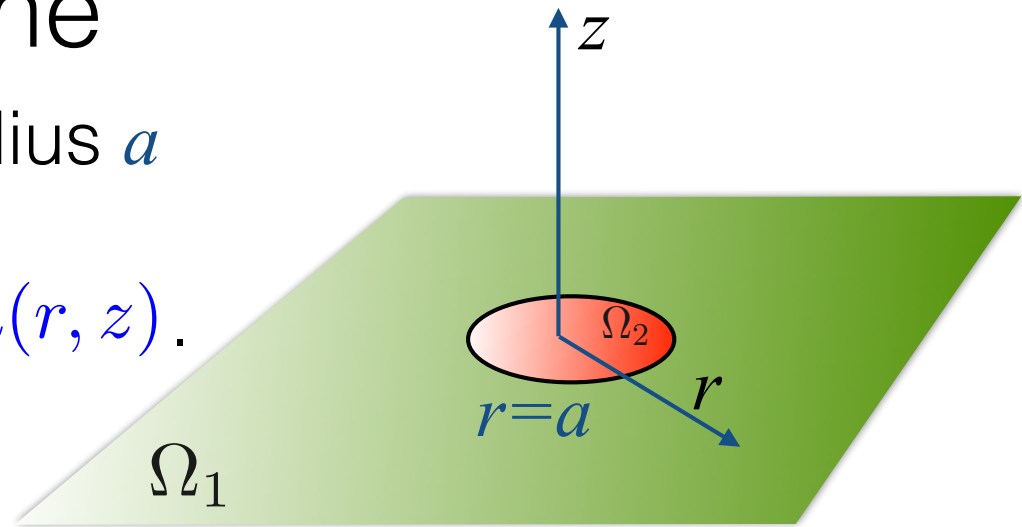
On the surface, the flux is:

$$-u_z(r, 0) = \begin{cases} 0 & r > a \\ \frac{2}{\pi} \frac{C_\infty}{\sqrt{a^2 - r^2}} & r < a \end{cases}$$

The total flux into the pore is:

$$\mathcal{F}_{pore} = 4aC_\infty$$

**Flux ~ Perimeter !!**





# Homogenization: Many pores on a plane

For  $N$  pores of radius  $a$ :

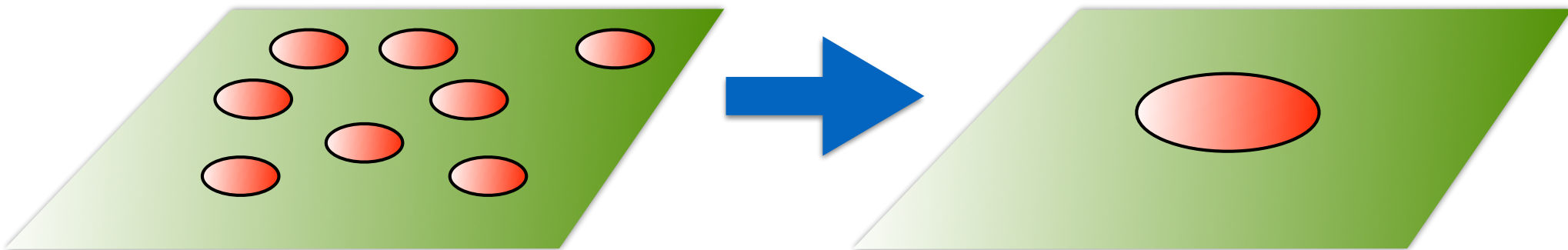
$$\mathcal{F}_{pore} = 4aN C_{\infty} + \frac{16a^2}{\pi} \sum_{i=1}^N \sum_{\substack{j=1 \\ j \neq i}}^N \frac{1}{d_{ij}}$$

Berg & Purcell 1977

Bernoff & Lindsay 2016

When the distance between pores,  $d_{ij} \gg a$ .

This tells us how to replace many pores with one.

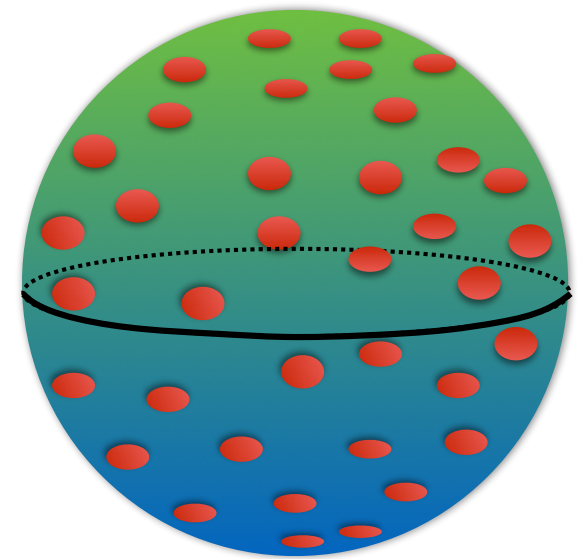


# Homogenization: Pores on a Sphere

This is a classic problem considered by Berg & Purcell (1977) in the dilute pore limit.

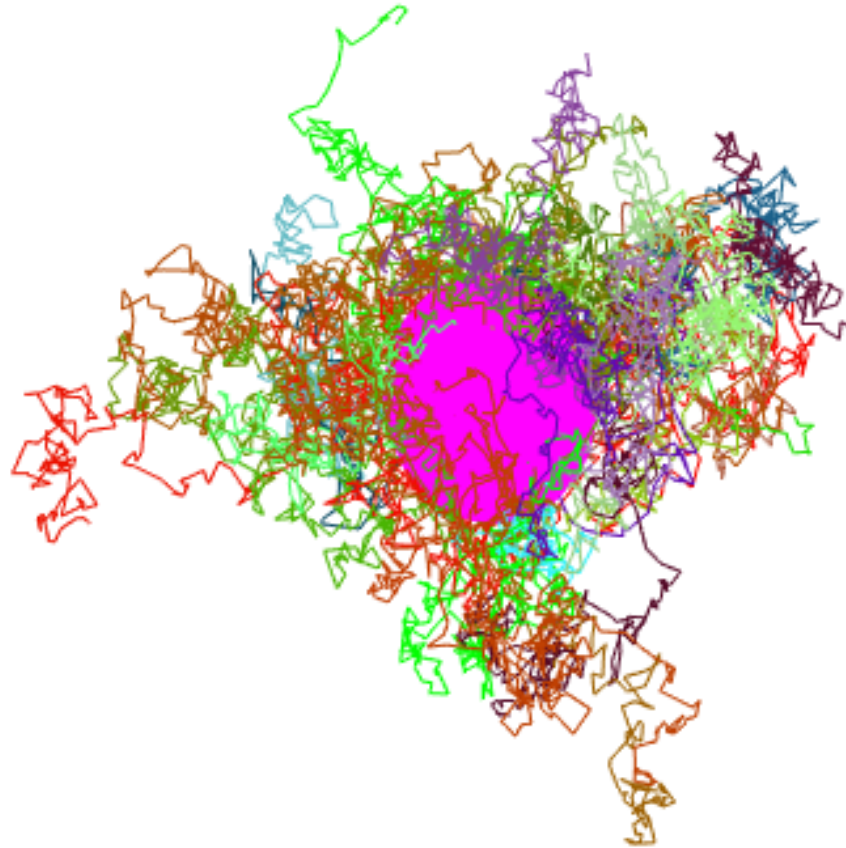
$$\mathcal{F}_{total} = 4aDNC_{\infty} \cdot \frac{1}{1 + Na/\pi R}$$

Knowing the flux allows you to replace the sphere with either a smaller sphere or a **Robin** BC.



The next order correction in the asymptotics has been computed by Lindsay & Ward (2016).

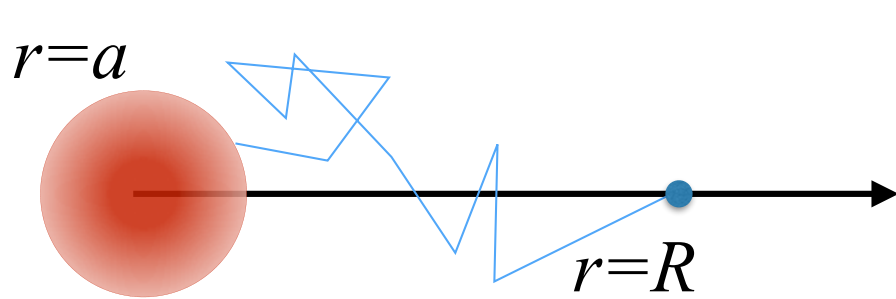
# Random Walk Numerics



random walk simulation for an absorbing sphere

# 3D Random Walk Numerics

Consider a  $\delta$ -function release at  $r=R$  and a spherical trap of radius  $a$ .



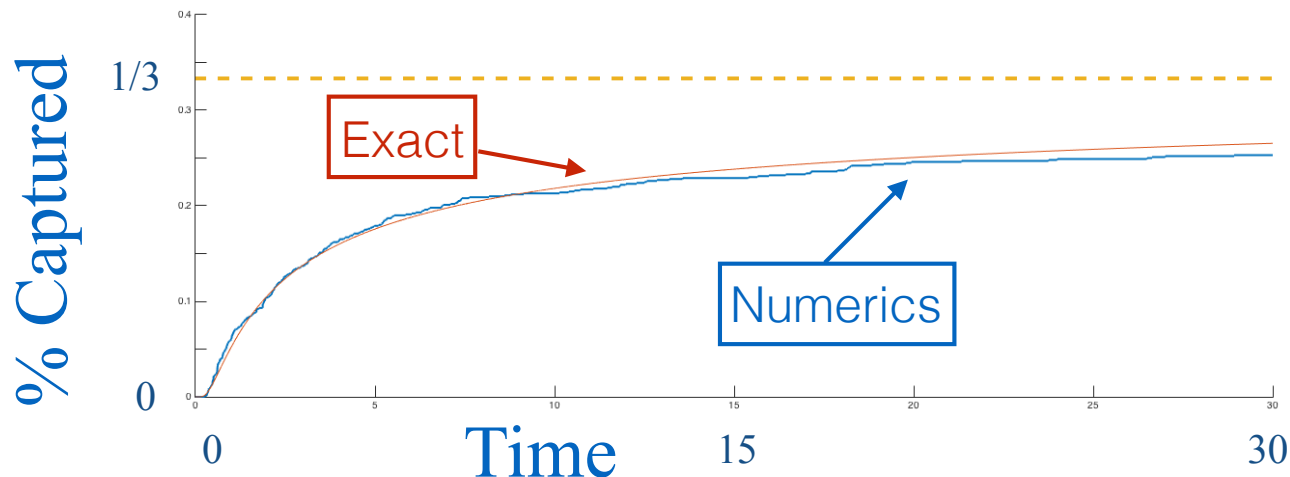
$$u_t = D \left[ u_{rr} + \frac{2}{r} u_r \right] \quad r > 0, t > 0$$
$$u(r, 0) = \frac{1}{4\pi R^2} \delta(r - R) \quad r > a$$
$$u(a, t) = 0 \quad t > 0$$

Numerical Method:

- $N$  particles
- Take Gaussian distributed steps of **Variance** =  $D\Delta t$
- Stop when particle is inside target.

$$a = 1, R = 3, D = 1$$

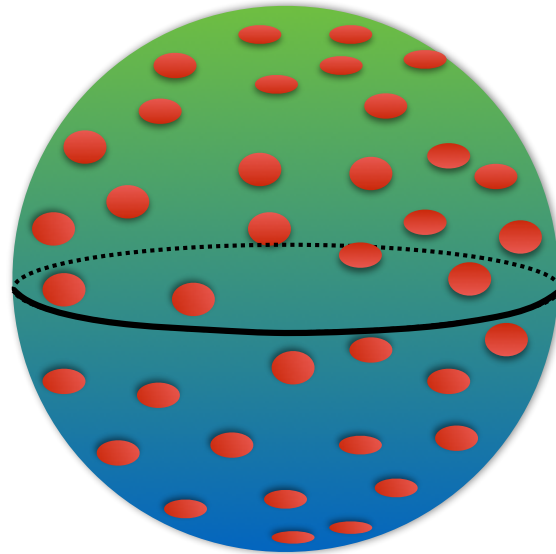
$$N = 1000, \Delta t = .001$$



Challenges:

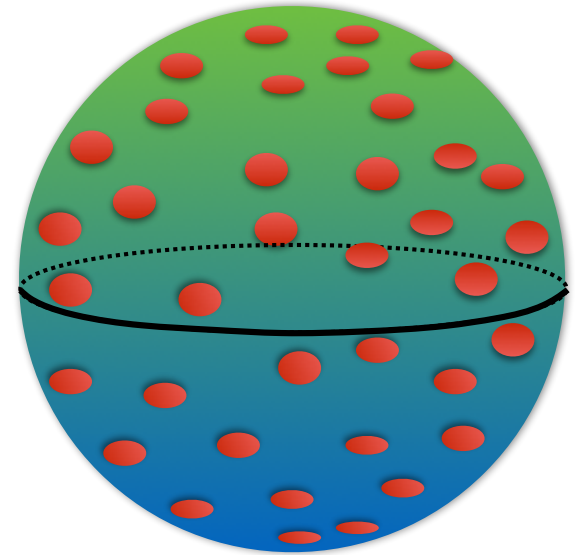
- Slow convergence.
- Long-time asymptotics often impossible to verify.
- Good for capture %, but difficult to obtain capture rate.

# Challenge Problems



# A Challenge Problem

Suppose a set of particles are released at the center of a sphere of radius  $R$ . The sphere has  $N$  pores of radius  $a$  on its surface. Compute the distribution of exit times for the sphere.



Notes:

- You can do this fairly effectively with a random walk particle code.

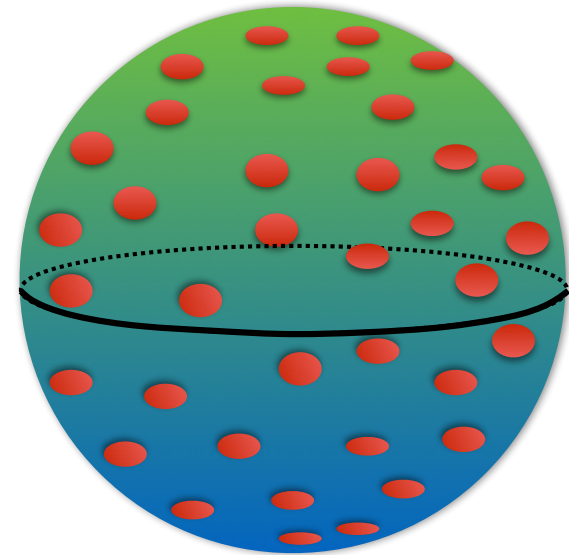
# The Homogenized Challenge Problem

Suppose a set of particles are released at the center of a sphere of radius  $R$ . The exit of the particles is modeled by a *Robin* boundary condition. Compute the distribution of exit times for the sphere.

$$u_t = D \left[ u_{rr} + \frac{2}{r} u_r \right] \quad R > r > 0, \quad t > 0$$

$$u(r, 0) = \frac{1}{4\pi r^2} \delta(r) \quad R > r > 0,$$

$$u_r(R, t) + \kappa u(R, t) = 0 \quad t > 0$$



Notes:

- You may wish to start out with the 1D version.
- Make sure you consider  $\kappa, D$  small and large.

# References

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- Bialek, William & Sima Setayeshgar. *Physical limits to biochemical signaling* *Proceedings of the National Academy of Sciences of the United States of America* 102.29 (2005): 10040-10045.
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- Lindsay, Alan E. & Michael J. Ward. *First Passage Statistics for the Capture of a Brownian Particle by a Structured Spherical Target with Multiple Surface Traps*. Preprint (2016).